



ESD & Shielding



Triboelectric effect

The **triboelectric effect** (also known as *triboelectric charging*) is a type of contact electrification in which certain materials become electrically charged after they come into contact with another different material through friction

Negative	Conductor - Insulator
Silicone rubber	Insulator
Polyethylene	Insulator
Saran	Insulator
Synthetic rubber	Insulator
Silver	Conductor
Nickel/copper	Conductor
Hard rubber	Insulator
Steel	Conductor
Cotton	Insulator
Paper	Insulator
Aluminum	Conductor
Silk	Insulator
Wool	Insulator
Nylon	Insulator
Glass	Insulator
Positive	Conductor - Insulator

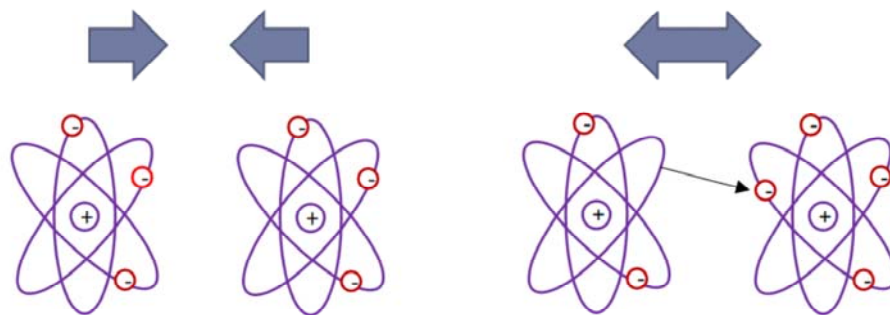
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Triboelectric effect

Electrostatic discharge produces:

- Currents of several amps
- High voltages (kV)
- $t < 100\text{ns}$
- Large bandwidth.

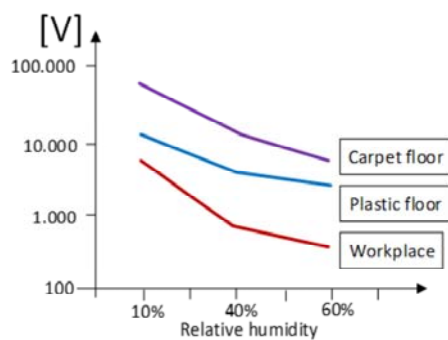


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The friction between the materials causes transfer of electrons and electrostatic charge of the materials



Electrostatic discharge



Typical voltage levels at different relative humidities

Generation	0-25% RH	65-90% RH
Walking across carpet	35.000 V	1.500 V
Walking across vinyl tile	12.000 V	250 V
Worker at bench	6.000 V	100 V
Chair with urethane foam	18.000 V	1.500 V



Electrostatic discharge

Insulators

A considerable amount of charge can be generated on the surface of an insulator. Because an insulating material does not readily allow the flow of electrons, both positive and negative charges can reside on an insulating surface at the same time, although at different locations

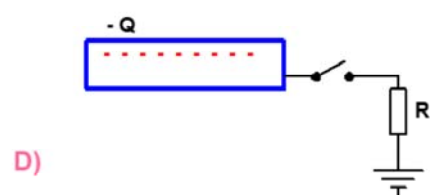
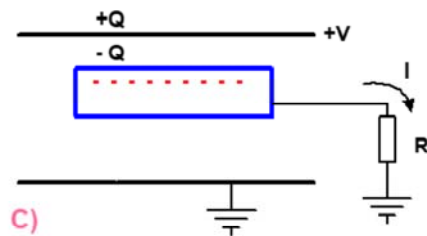
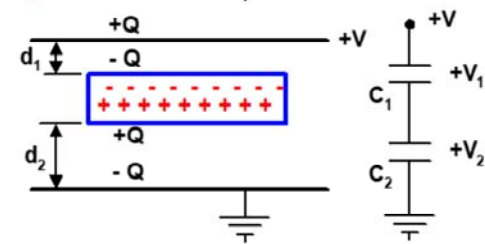
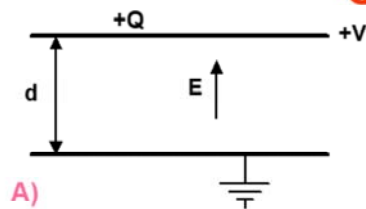
Conductive Materials

Due to its low electrical resistance, conductive materials allows electrons to flow easily across its surface or through its volume. When a conductive material becomes charged, the charge will be uniformly distributed across the surface of the material.

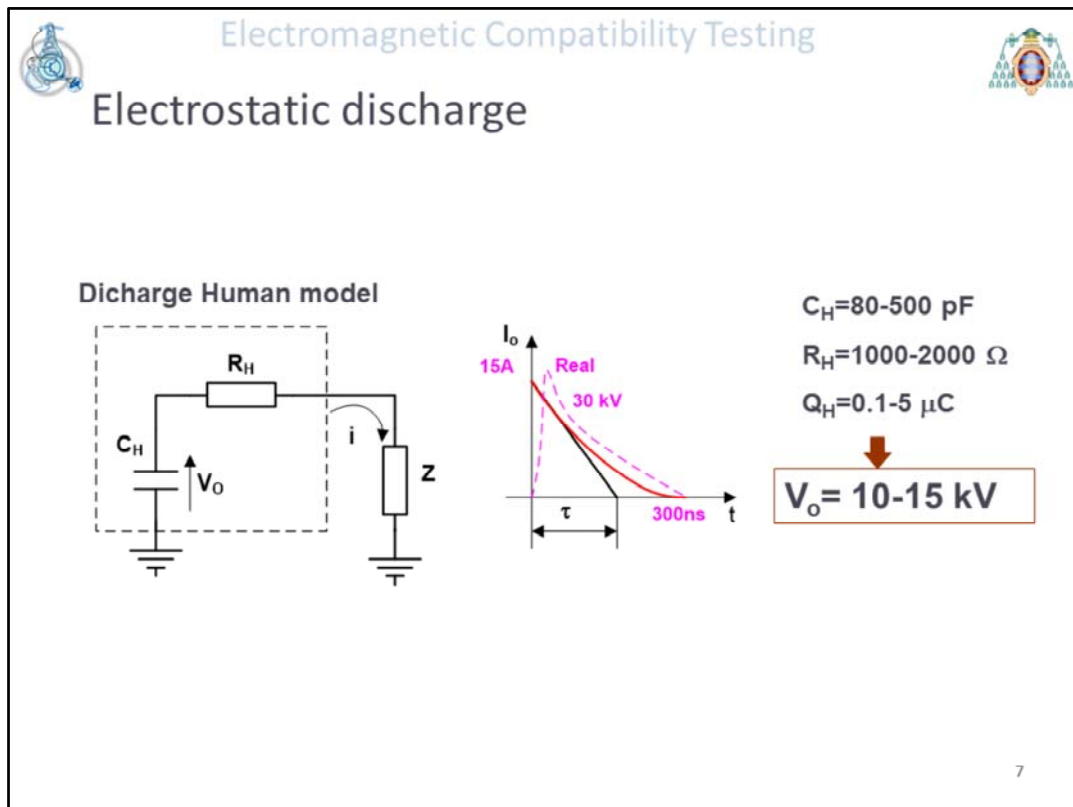


Charging process

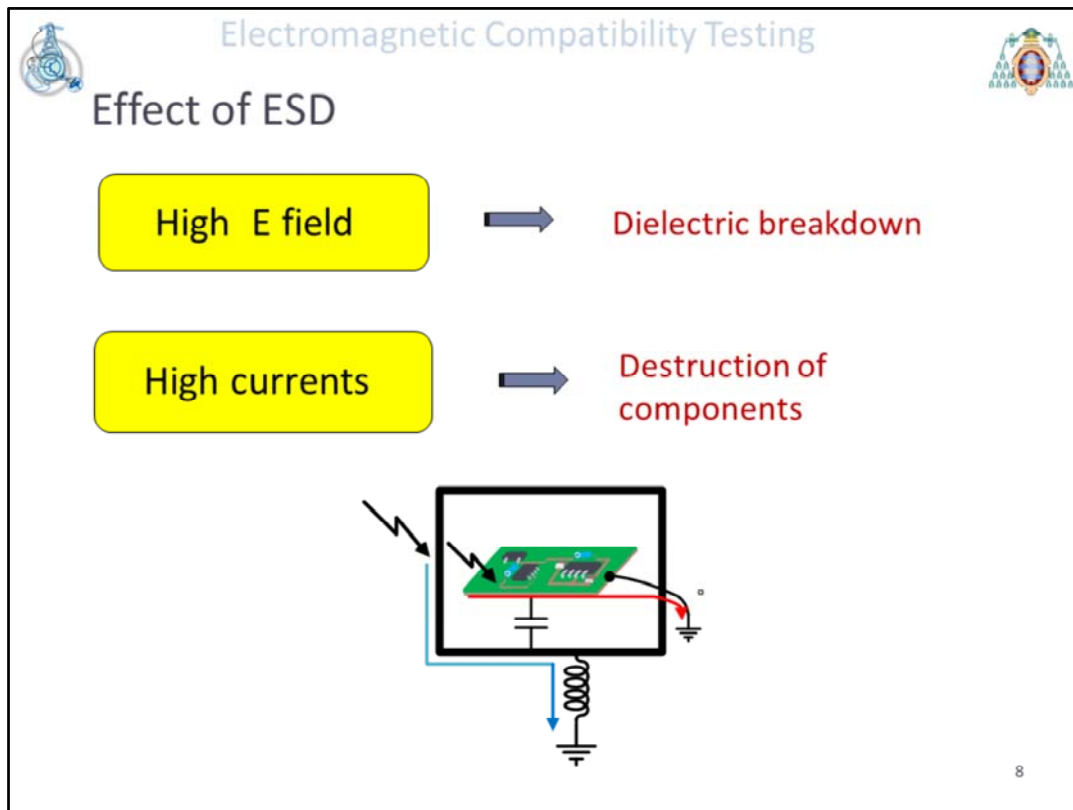
$$Q = C \cdot V$$



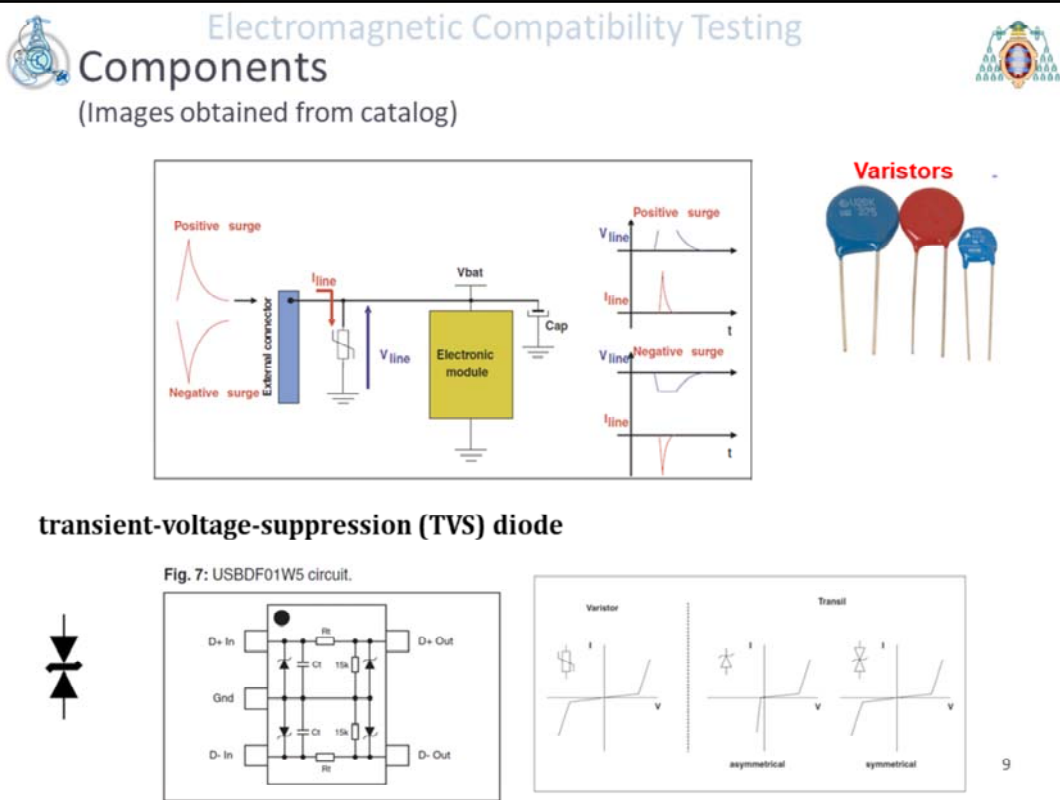
An external field cannot generate any field inside a conductive enclosure $\rightarrow$ Antistatic covers, paint, etc.



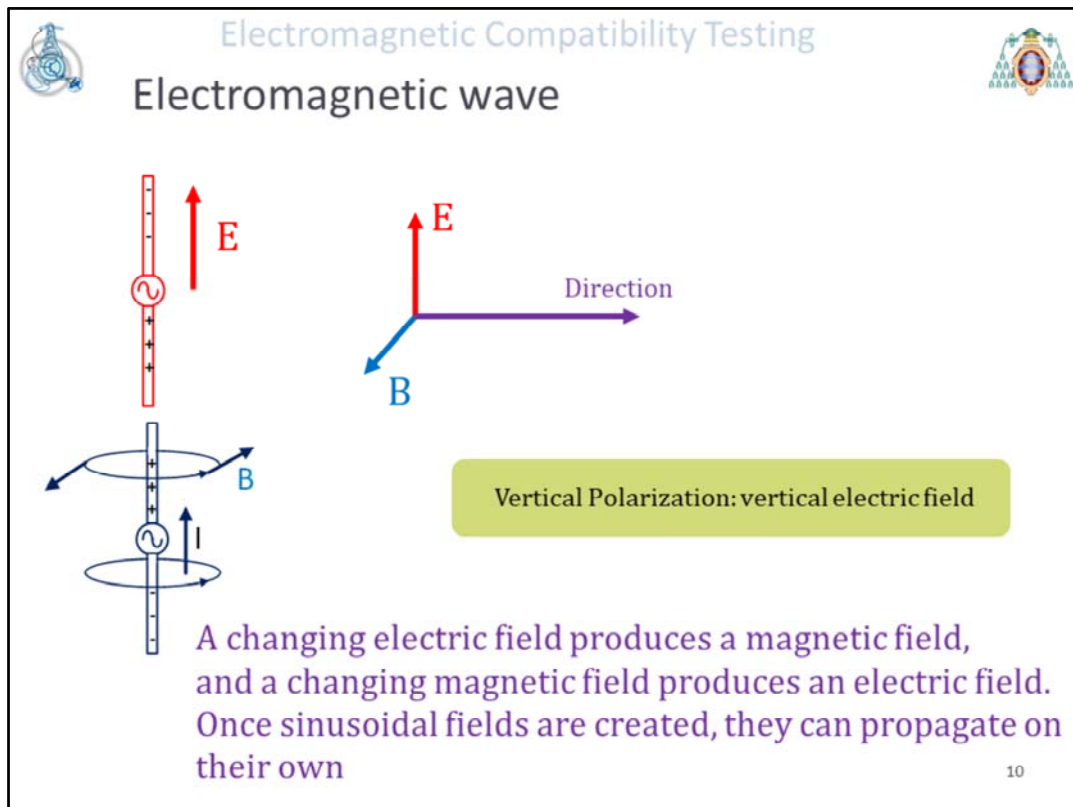
The process of electrostatic discharge caused by the human being through a resistance and a capacitor can be characterized with a model. The voltages reached can reach thousands of volts and tens of amps



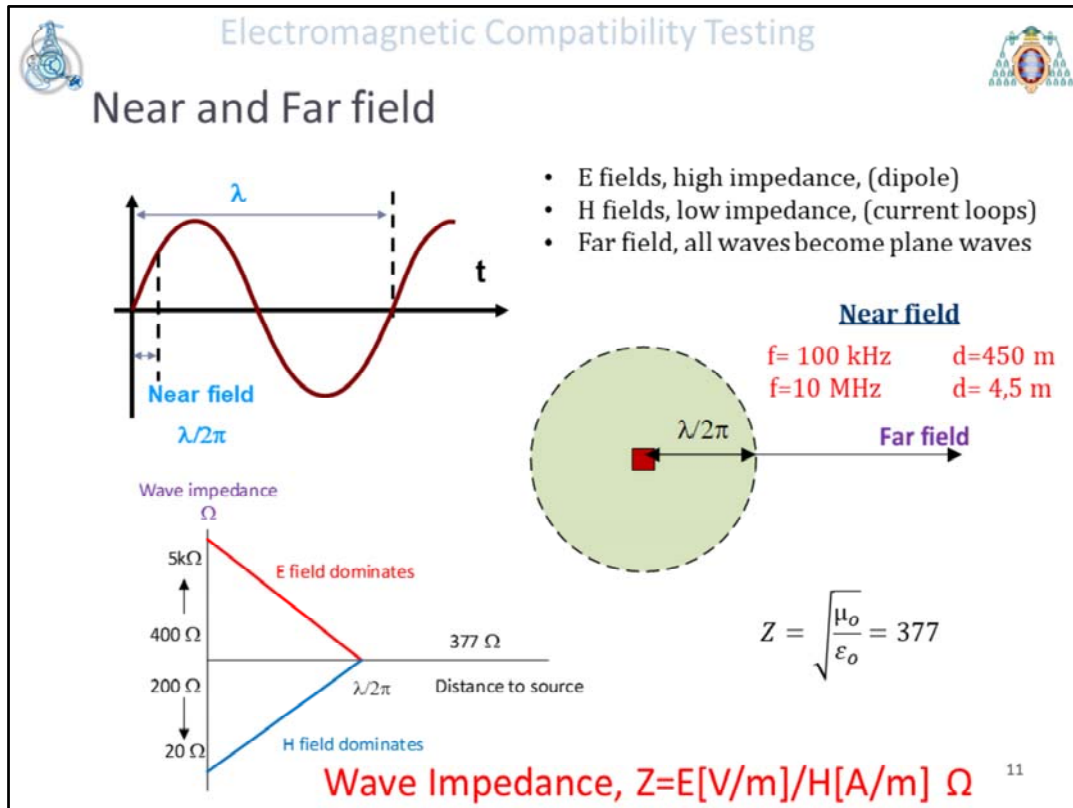
Encapsulating circuits in metallic boxes is not a solution to prevent electrostatic discharges if the ground connection is not adequate. Providing alternative paths to ground for the conduction of currents caused by ESD is a recommended practice



The components used to eliminate the effect of electrostatic discharges are varistors and transils (TVS). They have a sufficiently fast response to suppress overvoltages, fixing the nominal voltage value of the device.



The electromagnetic wave is formed by an electric field perpendicular to a magnetic field and both perpendicular to the direction of propagation



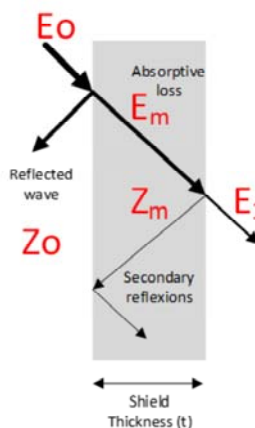
The term “near field” is used when the distance to the focus is less than the wavelength divided by two times pi



Electromagnetic Compatibility Testing



Electromagnetic waves



Incident wave: E_0 , Z_0

Reflected wave: E_m , Z_m

Transmitted wave: E_1

Absorptive loss

Secondary reflexions

Shield Thickness (t)

Shielding effectiveness (S) is a measure of how well an enclosure attenuates electromagnetic fields

$S = 20 \log (E_0/E_1) \text{ (dB)}$

$S = 20 \log (H_0/H_1) \text{ (dB)}$

Theoretical S of homogeneous material

- Reflective losses, R
- Absorption losses, A and
- Secondary reflective losses, B (ignore if $A > 8 \text{ dB}$)

$S = A + R + B$

For conductors, shield impedance:

If $\sigma \gg j\omega\epsilon$

$$Z = \sqrt{\frac{j\omega\mu}{\sigma}}$$

$Z_0 = 377$:free space impedance

$Z_m = E_s/H_s$: intrinsic material impedance

Effectiveness of a shield is associated to the ratio intensity of the incident wave to the wave that is capable of crossing the shield.



Electromagnetic waves

The amplitude of the reflected field must satisfy the relation,

If $t \rightarrow \infty$ infinite thickness

$$|E_{ref}| = |E_o| \cdot \Gamma_E$$

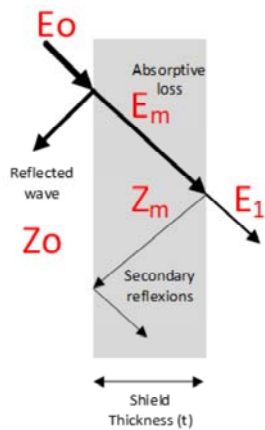
$$\Gamma_E = \frac{Z_m - Z_o}{Z_m + Z_o} \quad : \text{Field reflection coefficient}$$

$$|E_m| = |E_o| \cdot T_E$$

$$T_E = \frac{2Z_m}{Z_m + Z_o} \quad : \text{Field transmission coefficient}$$

If the material is lossy $\sigma \neq 0$, the transmitted wave decreases

$$|E_m(x)| = |E_m(x=0)| e^{-x/\delta} \quad \delta = \frac{1}{\sqrt{\pi f \mu \sigma}} \quad \text{Skin depth}$$

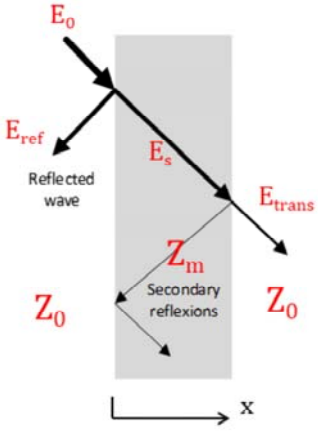




Electromagnetic Compatibility Testing

Electromagnetic waves



$$T_{E2} = \frac{2Z_0}{Z_m + Z_0}$$


Eo->m $|E_{trans}| = |E_m(x=t)| \cdot T_{E2}$

Then

$$|E_{trans}| = |E_o| \left(\frac{2Z_m}{Z_0 + Z_m} \right) \left(\frac{2Z_0}{Z_0 + Z_m} \right) e^{-t/\delta}$$

$$Z = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}} \approx \sqrt{\frac{j\omega\mu}{\sigma}} \approx \sqrt{\frac{\omega\mu}{\sigma}} e^{j\pi/4}$$

This expression applies to any shield material that is much thicker than one skin depth

$$|E_{trans}| = |E_o| \cdot \frac{4Z_m}{Z_0} \cdot e^{-t/\delta}$$

$$S = 20 \log \frac{|E_o|}{|E_{trans}|} = \underbrace{20 \log \frac{Z_0}{4Z_m}}_{R(dB)} + \underbrace{20 \log e^{t/\delta}}_{A(dB)}$$

The final result shows the two most important terms: losses by reflection and losses by absorption.



Practical considerations

Far field (plane wave)

Low frequency → Reflective losses (use conductive material)
High frequency → Absorption losses (use magnetic material)

Near field

Electric field E

Low frequency → Reflective losses
High frequency → Absorption losses

Magnetic field H

Low and High frequency → Absorption losses



Electromagnetic Compatibility Testing

Absorption losses "A" (E, H and planar waves)



The absorption loss, $A(dB)$, is the attenuation due to power converted to heat

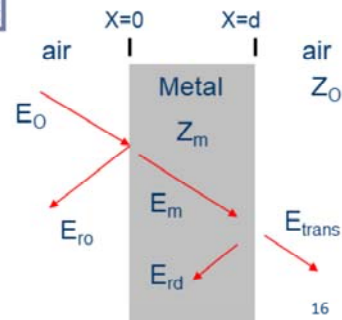
$$A = \frac{E_o}{E_d} = \frac{H_o}{H_d} = e^{d/\delta} = 20 \log e^{\frac{d}{\delta}} [dB] \approx 8.7 \cdot \frac{d}{\delta} [dB]$$
$$\delta = \sqrt{\frac{\rho}{\pi \mu f}}$$

A = attenuation dB
 d = Thickness
 f = frequency
 δ = skin depth
 ρ = resistivity
 H_d = field at $x=d$

H at low frequency Absorption losses
are the most important

8dB by each δ that increases d

Doubling the thickness of the shield doubles
the loss in decibels



Absorption losses depend fundamentally on thickness and skin depth.



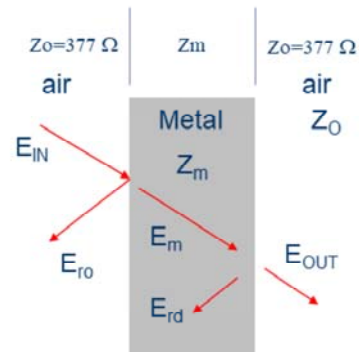
Reflective losses "R"

$$R = \frac{E_{IN}}{E_{OUT}} = \frac{H_{IN}}{H_{OUT}} = \frac{(Z_O + Z_m)^2}{4(Z_O Z_m)}$$

$$Z_O = 377 \Omega$$

$$|Z_m| = \frac{E_M}{H_M} = \sqrt{\left| \frac{j\omega\mu}{j\omega\epsilon + \sigma} \right|} = \frac{\sqrt{2}}{\sigma \cdot \delta}$$

$$\delta = \sqrt{\frac{\rho}{\pi\mu f}}$$



Reflective losses depend on the impedances of the different media involved

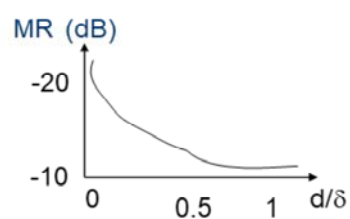
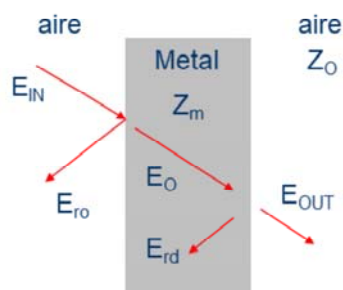


Secondary reflective losses "MR"

$$MR = 20 \cdot \log \left(1 - e^{-\frac{2d}{\delta}} \right) = \frac{2d}{\delta} [dB]$$

(for low values of d/δ)

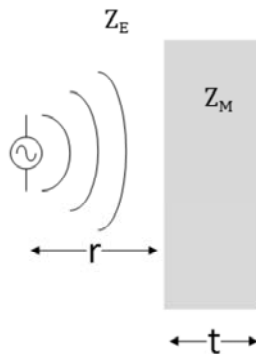
MR is important if absorption losses are low





Near field: E dipole

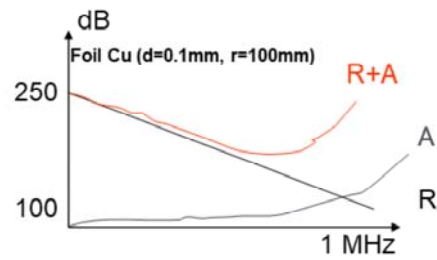
$$Z_E = \frac{1}{2\pi \cdot \epsilon_o \cdot f \cdot r}$$



r: distance from the source to the shield

Characteristics

- High conductivity
- Zero voltage or constant voltage

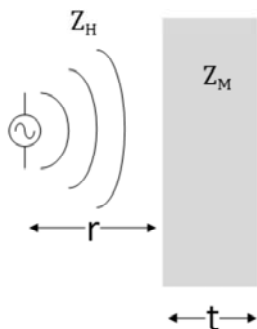


$$R_E = \left| \frac{(Z_E + Z_m)^2}{4 Z_E Z_m} \right| \approx \frac{1}{4\sqrt{2}} \cdot \frac{\sigma \delta}{w \cdot \epsilon_o \cdot r}$$

If $w \rightarrow 0 \Rightarrow R_E$ will be high₁₉



Near field: H dipole



$$Z_H = 2 \pi \mu f \cdot r$$

$$R_H = \left| \frac{(Z_H + Z_m)^2}{4 Z_H Z_m} \right| \approx \frac{w \mu_o r \sigma \delta}{4 \sqrt{2}}$$

If $Z_H \gg Z_m$. This equation is not correct

- Low freq. Low effectiveness.
- High freq. Absorción

$d=r$: distance from the source to the shield

- A_H (absorción) increases if f or r increase
- A_H increases with high permeability
- Increase with "thickness"

magnetic materials are the most appropriate ones



Planar waves

$$Z_0 = 377 \Omega$$

Far field ($r \gg \lambda/2\pi$)

It is not very common

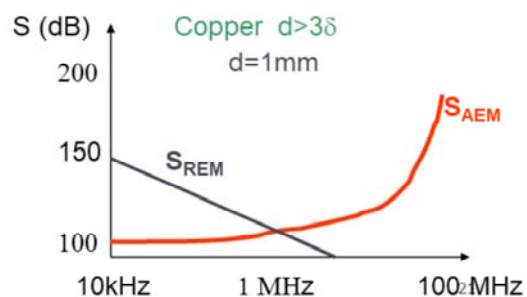
$$R_{EM} = \left| \frac{Z_o}{4 \cdot Z_m} \right| = \frac{377 \sigma \delta}{4 \sqrt{2}} \quad \text{Reflective losses}$$

At low frequency (d/δ low) MR and absorption are negligible

At $f \downarrow$ reflective losses are important.

- Good conductors (high σ)

$$\sigma \gg \omega \epsilon$$

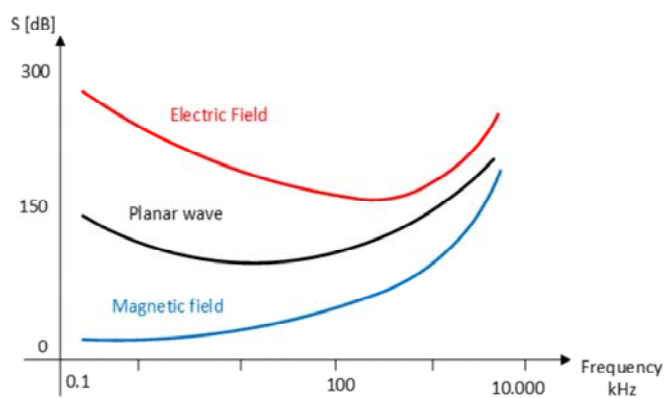




Electromagnetic Compatibility Testing



0.02-in-thick solid aluminium shield

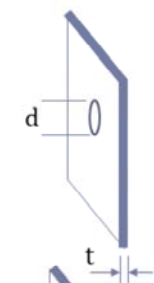


Electric field, plane wave, and magnetic field shielding effectiveness of a 0.02-in-thick solid aluminium shield.



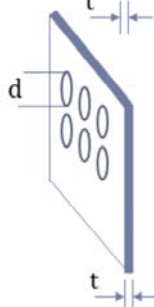
Practical considerations

Holes/Apertures $d > t$



$$\text{If } \lambda/2 < d \quad SE \approx 0 \text{ dB}$$

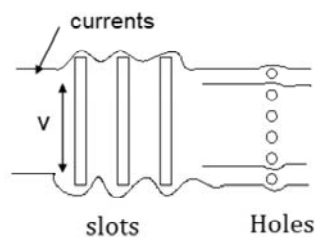
$$\text{If } \lambda/2 > d \quad SE \approx 20 \log (\lambda/2d)$$



$$\text{If } \lambda/2 < d \quad SE \approx 0 \text{ dB}$$

$$\text{If } \lambda/2 > d \quad SE \approx 20 \log (\lambda/2d) - 10 \log (n)$$

n: number of holes



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Waveguide

Circular waveguide

Cutoff frequency below which it becomes an attenuator



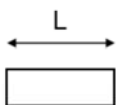
$$f_c = \frac{6.9 \cdot 10^9}{D} [Hz]$$

Effectiveness for $f < f_c$

$$\lambda_{\max} = \frac{D}{2\pi}$$

$$S = 32 \cdot \frac{t}{D} [dB]$$

Rectangular waveguide

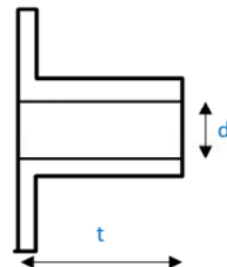


$$f_c = \frac{5.9 \cdot 10^9}{L} [Hz]$$

$$S = 27 \cdot \frac{t}{L} [dB]$$

$$\lambda_{\max} = \frac{L}{2\pi}$$

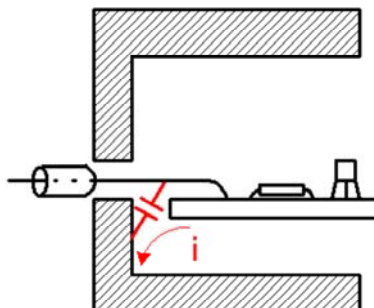
* Length in inches
t: thickness



When the thickness is similar to the hole dimension



Electromagnetic Compatibility Testing





Magnetic shielding

**WE-FAS Flexible
Absorber Sheet**



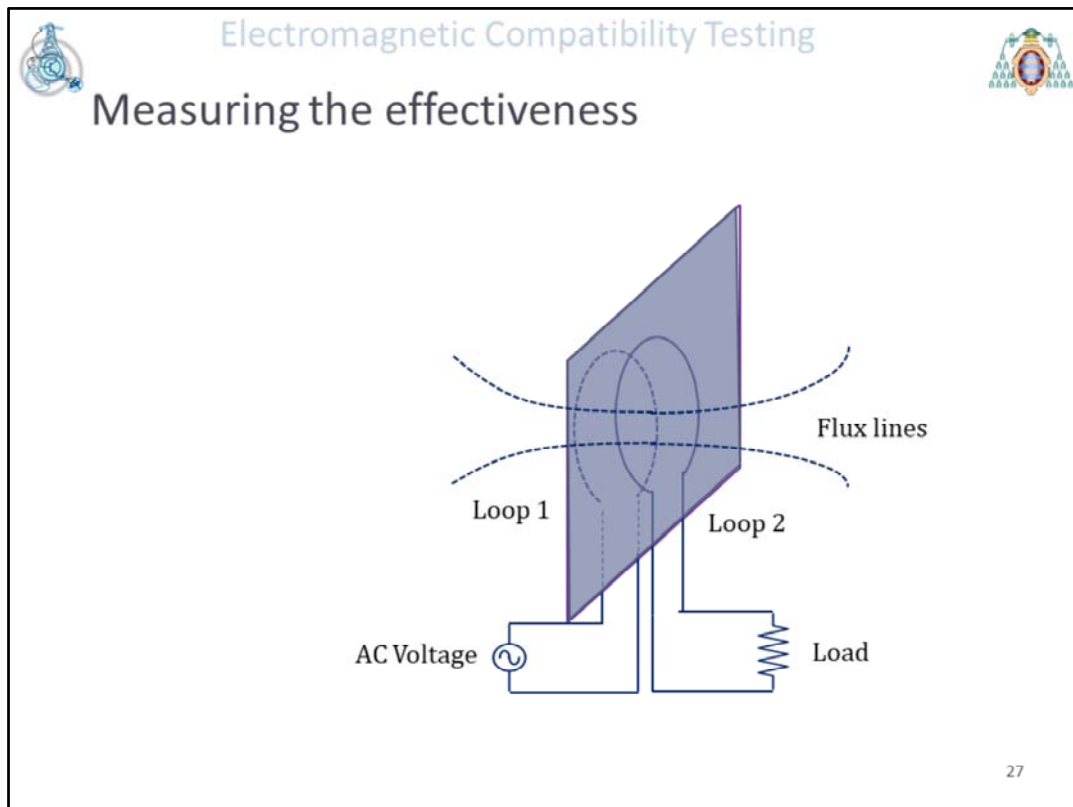
**WE-FSFS Flexible
Sintered Ferrite
Sheet**



Images obtained from Wurth Electronics catalogue

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There are new magnetic adhesive materials to shield small components



To measure the effectiveness of a shielding you can use two loops, one connected to a source and the other one connected to the load. Placing the shield between both loops we will observe the voltage that appears across the load due to the field that crosses the shield